

MQ I Particule chargée dans un champ  
magnétique uniforme

$$\vec{B} = \text{rot } \vec{A} = \vec{\nabla} \times \vec{A}$$

$$H = \frac{(\vec{p} - q\vec{A})^2}{2m}$$

$$1) [\vec{p}, \vec{A}] = \vec{p}\vec{A} - \vec{A}\vec{p}, \quad \vec{p} = -i\hbar\vec{\nabla}$$

$$\begin{aligned} \circ [\vec{p}, \vec{A}] f &= (-i\hbar\vec{\nabla})(\vec{A} f) - (\vec{A}(-i\hbar\vec{\nabla})f) \\ &= -i\hbar(\vec{\nabla}\vec{A})f - \cancel{i\hbar\vec{A}\vec{\nabla}f} + \cancel{i\hbar\vec{A}\vec{\nabla}f} = -i\hbar \text{div}(\vec{A}) f \end{aligned}$$

$$\Rightarrow [\vec{p}, \vec{A}] = -i\hbar \text{div} \vec{A}$$

$$\Rightarrow \text{jauge de Coulomb: } \text{div} \vec{A} = 0 \Rightarrow [\vec{p}, \vec{A}] = 0$$

$$\Rightarrow \vec{p}\vec{A} = \vec{A}\vec{p}$$

$$\Rightarrow H = \frac{1}{2m} (\vec{p} - q\vec{A})(\vec{p} - q\vec{A})$$

$$= \frac{1}{2m} (\vec{p}^2 - q\vec{p}\vec{A} - q\vec{A}\vec{p} + q^2\vec{A}^2)$$

$$= \frac{1}{2m} (\vec{p}^2 - 2q\vec{A}\vec{p} + q^2\vec{A}^2)$$

$$2) \vec{A} = Bx\vec{\mathbb{1}}_y, \quad \vec{B} = \text{rot} \vec{A}$$

$$\vec{\nabla} \times \vec{A} = \begin{pmatrix} \partial_x \\ \partial_y \\ \partial_z \end{pmatrix} \times \begin{pmatrix} 0 \\ Bx \\ 0 \end{pmatrix} = \begin{pmatrix} \partial_y \cdot 0 - \partial_z \cdot Bx \\ \partial_z \cdot 0 - \partial_x \cdot 0 \\ \partial_x \cdot Bx - \partial_y \cdot 0 \end{pmatrix} = B\vec{\mathbb{1}}_z \rightarrow \text{uniforme} \checkmark$$

$$\text{div } \vec{A} = 0 ?$$

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$$\begin{aligned}\vec{\nabla} \cdot \vec{A} &= \partial_x A_x + \partial_y A_y + \partial_z A_z \\ &= \partial_x \cdot 0 + \partial_y (Bx) + \partial_z \cdot 0 = 0 \checkmark\end{aligned}$$

$$H = \frac{1}{2m} (\vec{p}^2 - 2q \vec{A} \cdot \vec{p} + q^2 \vec{A}^2)$$

$$\vec{A} \cdot \vec{p} = Bx p_y, \quad \vec{A}^2 = B^2 x^2$$

$$\Leftrightarrow H = \frac{1}{2m} \left( -\frac{\hbar^2}{2m} \Delta + i\hbar 2qBx \partial_y + q^2 B^2 x^2 \right)$$

3)

$$[H, p_x] = \frac{1}{2m} \left( -\frac{\hbar^2}{2m} [\vec{p}^2, p_x] - 2qB [x p_y, p_x] + q^2 B^2 [x^2, p_x] \right)$$

$$= \frac{1}{2m} \left( -2qB \left( x [p_y, p_x] + [x, p_x] p_y \right) + q^2 B^2 2ix \right)$$

$$= i\hbar \frac{qB}{m} (-p_y + qBx) \neq 0$$

$$[H, p_y] = \frac{1}{2m} \left( -\frac{\hbar^2}{2m} [\vec{p}^2, p_y] - 2qB [x p_y, p_y] + q^2 B^2 [x^2, p_y] \right) = 0$$

$$[H, p_z] = \dots = 0 \quad \Rightarrow \text{E.C.O.C. } H, p_y, p_z \quad \text{car } [p_y, p_z] = 0$$

(ensemble complet d'observables qui commutent)

M&I

$$4) \psi(x, y, z) = \Phi(x) e^{ik_y y} e^{ik_z z}$$

$$H\psi = E\psi$$

$$\frac{1}{2m} (-\hbar^2 \Delta + i\hbar 2qBx \partial_y + q^2 B^2 x^2) \Phi(x) e^{ik_y y} e^{ik_z z} = E\psi(x, y, z)$$

$$= \frac{1}{2m} \left( -\hbar^2 \Phi''(x) e^{ik_y y} e^{ik_z z} - \hbar^2 \Phi(x) (-k_y^2 e^{ik_y y}) e^{ik_z z} \right.$$

$$\left. - \hbar^2 \Phi(x) (-k_z^2) e^{ik_y y} e^{ik_z z} + i\hbar 2iqk_y Bx \Phi(x) e^{ik_y y} e^{ik_z z} + q^2 B^2 x^2 \Phi(x) e^{ik_y y} e^{ik_z z} \right)$$

$$= \left( -\frac{\hbar^2}{2m} \Phi''(x) + \frac{\hbar^2}{2m} (k_y^2 + k_z^2) \Phi(x) - \frac{\hbar q k_y Bx}{m} \Phi(x) \right.$$

$$\left. + \frac{q^2 B^2 x^2}{2m} \Phi(x) \right) e^{ik_y y} e^{ik_z z} = E\psi(x, y, z)$$

$$\frac{-\hbar^2}{2m} \Phi''(x) + \frac{q^2 B^2}{2m} (x - x_0)^2 \Phi(x) = E_{n_x} \Phi(x) \Rightarrow \text{Oscillateur harmonique en } x$$

"shifté" par  $x_0$

$$5) x_0 = \frac{\hbar k_y}{Bq} = \frac{\hbar q}{qB}$$

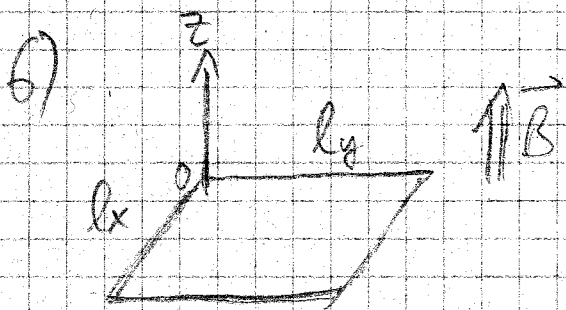
$$E_{n_x} = E - \frac{\hbar^2 k_z^2}{2m} \Rightarrow E_{n_x} = \left(n + \frac{1}{2}\right) \hbar \omega \quad (n = n_x)$$

$$\omega = \frac{qB}{m} = \text{fréquence cyclotron}$$

$$E = E_{n_k z k_y} = \left(n + \frac{1}{2}\right) \frac{\hbar g B}{m \omega_c} + \frac{\hbar^2 k_z^2}{2m}$$

Fonction d'ondes pour l'oscillateur harmonique donner par les polynômes d'Hermite

$$\Rightarrow \Psi_{n_k y k_z}(x, y, z) \propto e^{i(k_z z + k_y y)} e^{-\left(x - \frac{\hbar k_y}{g B}\right)^2 \frac{g B}{2 \hbar}} H_n \left( \sqrt{\frac{g B}{\hbar}} \left(x - \frac{\hbar k_y}{g B}\right) \right)$$



Pas de mouvement en  $z \Rightarrow k_z = 0$

Conditions aux bords:  $\Psi(y=0) = \Psi(y=l_y) = 0$

$$\Rightarrow \Psi \propto \sin(k_y y) \text{ avec } k_y = \frac{n_y \pi}{l_y}, n_y \in \mathbb{Z}$$

Dégénérescence:  $0 < x_0 \leq l_x$

$$\Leftrightarrow 0 < \frac{\hbar k_y}{g B} \leq l_x$$

$$\Leftrightarrow 0 < n_y \leq \frac{B l_x l_y}{\frac{\hbar^2}{g}} = \frac{\Phi_B}{\Phi_0} \leftarrow \text{quantum de flux magnétique}$$

Flux magnétique

pour  $B = 2,60 \text{ T}$ ,  $l_x = l_y = 10^{-2} \text{ m}$

(longueur caractéristique puits:  $\sqrt{\frac{\hbar}{g B}} \approx 2,57 \cdot 10^{-2} \text{ m} \ll l_x, l_y$ )

$$\Rightarrow n_y \leq 2,0048 \cdot 10^{10}$$

⇒ Mesurer  $\omega$  (par absorption de microondes) permet de mesurer la masse effective des électrons dans un matériau

Séance 6-5

laser  $\lambda = 337 \mu\text{m}$ , raie abs. pour  $B = 2,605 \text{ T}$

$$\frac{qB}{m_f} = \omega = \frac{c2\pi}{\lambda} \Rightarrow \underline{\underline{\text{masse effective } m_f \approx 0,080 m_e}}$$

7)

Changement de jauge:

$$\vec{A} \rightarrow \vec{\tilde{A}} = \vec{A} + \vec{\nabla}f$$

$$\Rightarrow \psi \rightarrow \tilde{\psi} = e^{i\omega(\vec{x})} \psi$$

↑  
phase

Montrer que  $\langle \psi | (\hat{p} - q\vec{A}) | \psi \rangle = \langle \tilde{\psi} | (\hat{p} - q\vec{\tilde{A}}) | \tilde{\psi} \rangle$

(ça implique  $\langle \psi | \frac{(\hat{p} - q\vec{A})^2}{2m} | \psi \rangle = \langle \tilde{\psi} | \frac{(\hat{p} - q\vec{\tilde{A}})^2}{2m} | \tilde{\psi} \rangle$ )

$$\langle \tilde{\psi} | (\hat{p} - q\vec{\tilde{A}}) | \tilde{\psi} \rangle$$

$$= \langle \psi | e^{-i\omega} \left[ \underbrace{-i\hbar \vec{\nabla}}_{=\hat{p}} (e^{i\omega} \psi) + e^{i\omega} (-i\hbar \vec{\nabla}) \psi - q\vec{A} e^{i\omega} \psi - q(\vec{\nabla}f) e^{i\omega} \psi \right]$$

$$= \langle \psi | e^{-i\omega} \left( -i\hbar (\vec{\nabla} e^{i\omega}) - q(\vec{\nabla}f) e^{i\omega} \right) | \psi \rangle + \langle \psi | (\hat{p} - q\vec{A}) | \psi \rangle$$

$$\stackrel{!}{=} \langle \psi | e^{-i\omega} \left( -i\hbar \cdot i \vec{\nabla} \omega e^{i\omega} - q \nabla f e^{i\omega} \right) | \psi \rangle \stackrel{!}{=} 0$$

$$\Rightarrow \frac{1}{\hbar} \nabla \omega = q \nabla f \Rightarrow \boxed{\omega(\vec{x}) = \frac{qf(\vec{x})}{\hbar}}$$

$$\Rightarrow \boxed{\tilde{\psi} = e^{iqf(\vec{x})/\hbar} \psi}$$

Ex:  $\vec{\tilde{A}} = -B_y \mathbb{1}_x = \vec{A} + \vec{\nabla}f \Rightarrow f(\vec{x}) = Bxy$

$$\Rightarrow \tilde{\psi} = e^{iqBxy/\hbar} \psi$$